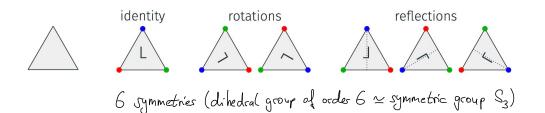
Ulrich Thiel https://ulthiel.com/math

What are "Lie Algubras" about ?

First, a clarification:

(pictures from wikipedia, public domain) Bruce Lee (1940-1973) Notice the difference in spelling! Juess you remember what a group is. I'm sure you also know where this concept comes from: Symmetry

E.g. every polynomial has a Galois group permuting the roots. A geometric abject has a group of symmetries leaving the object unchanged, e.g. for a triangle in the plane



(picture from own Tikz)

Knowing about symmetries is very important as it helps to simplify things. We mathematicians can study symmetry without necessarily having a concrete object/problem in mind. Can do this by studying an (abstract) group G. This brings us to the following concept.

A (linear) representation of a group G is a group morphism $g:G \longrightarrow GL(V)$ for some vector space V. Note that since g is a group morphism, every relation that holds in G, also holds between the linear operators $g(g), g\in G$, on V, i.e. the symmetry encoded abstractly in G is now more concretely also present in these linear operators. So, basically, a representation is a linear object obeying the symmetries encoded in G.

You know what a subspace $U \subseteq V$ is, right? So, what's a <u>suprepresentation</u> of jo?It's a subspace $U \subseteq V$ which is stable under all the operators p(g), $g \in G$. Here's an example: consider the representation

$$g: S_3 \longrightarrow GL_3(\mathbb{R}) = GL(\mathbb{R}^3)$$

$$\sigma \longmapsto (e_i \mapsto e_{\sigma(i)})$$

$$\uparrow_{i-th slandard basis kect}$$

Let $U = R^3$ be the line spanned by $e_1 + e_2 + e_3$. This is stable under S_3 , i.e., it is a subrepresentation. A representation is called irreducible if it contains no non-trivial (i.e. = 0, V) subrepresentation. Little exercise: in the S3 example above, convince yourself that we get an induced representation on R3/U, and that this representation is irreducible. What's the point of this? It is a classical fact (Maschke's theorem) that if G is finite and we work with representations over a held of characteristic zero (Q, R, C, ...), then every representation is (uniquely) a direct sum of irreducible representations. Hence, if we know all irreducible representations, we know all representations. => irreducible representations are the building blocks of (finite, linear) symmetry. Sz for example has 3 irreducible representations: the 2-dimensional we found above and two further 1-dimensional ones (which?). In an abstract group, the elements can't see each other: g h "I'm so lonely" "I'm so lonely" But actually, very often the symmetry transformations form some "geometric space" and can "see" each other. Consider for example the orientation preserving symmetries of the 2-sphere in R3. This is the special orthogonal group SO(3). It contains for example

$$R_{2}(\Theta) := \begin{pmatrix} \cos \Theta & -\sin \Theta & O \\ \sin \theta & \cos \Theta & O \\ O & O & 1 \end{pmatrix}, \quad \Theta \in [0, 2\pi]$$

around the z-axis. They vary "continuously" with a parameter O and we can think of R2(0) for "small" O as an "infinitesimal generator" Lz of rotations around the z-axis. This infinitesimal generator is a "tansent vector" to the "space" SO(3).

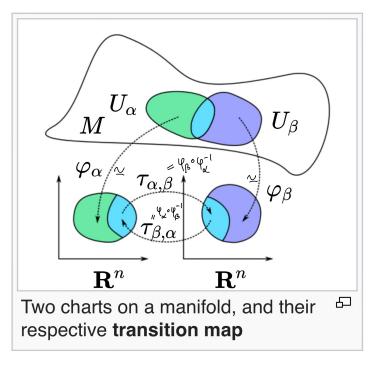
But how is this a space ? What is space?

the rotations

Our universe is probably not flat (like RM) but curved, e.g.



But <u>locally</u> (in a small neighbourhood of any point) it is flat, i.e. looks like (an open subset of) Rⁿ. The precise notion of a space which "(avally looks like Rⁿ" is that of a <u>manifold</u>. This is a (topological) space M which is glued together from open subsets of Rⁿ:



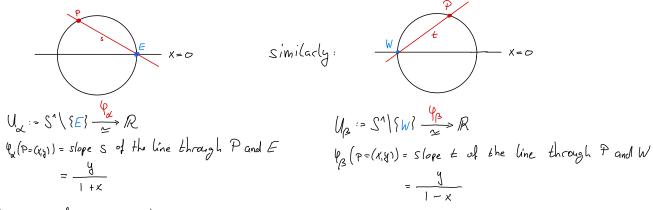


Browser screenshot, check out

quantamagazine.org/whatis-the-geometry-of-theuniverse-20200316/

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(picture from Wikipedia,
author Stomatipoll, Licene CC BK-SA 3.0)
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Here's an example: the circle $S^1 = \{x \in \mathbb{R}^2 \mid \|x\|^2 = 1\}$. How is this a manifold? It is patched together from two parts:

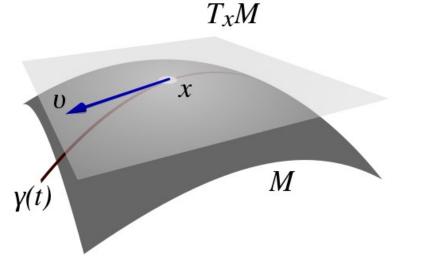


Transition function on the overlap is $T_{A} = (p \circ p^{-1} : R | So] \rightarrow R | So], S \mapsto \frac{1}{t}$. This is a diffeomorphism (differentiable isomorphism). => S¹ is a 1-dimensional manifold.

You can do a similar construction for the 2-sphere S², and more generally for S¹.

Let
$$M$$
 be a manifold. We want to define tangent directions in a point XEM.
Imagine, you sit in a car, drive with full speed on M , you drive through X, and then
exactly in X, your car gets out of conbrol and flies straight alf M . That's a tangent vector.
More precisely: take a chart (U, ψ) around X , $\psi: U \xrightarrow{\sim} V = \mathbb{R}^n$. Take a path
 $\chi: (-1, 1) \longrightarrow U \in X$

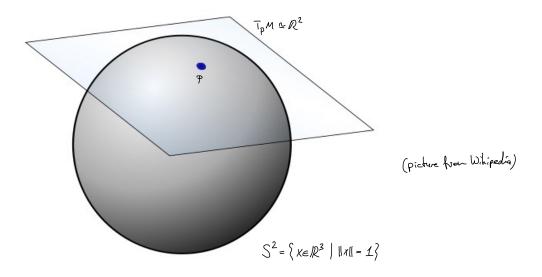
such that $\gamma(0) = x$ and such that $(\varphi \circ \chi : (-1, 1) \longrightarrow \mathbb{R}^n)$ is a differentiable map. Then we can think of $D((\varphi \circ \chi)(0) \in \mathbb{R}^n)$ as a tangent vector in x on M.



(picture from Wikipedia)

Two distinct paths can gield the same tangent rector \sim introduce an equivalence relation on paths \sim get a vector space $T_x M$ of equivalence classer $(\frac{\tan gent}{\tan gent}, \frac{\pi}{\tan gent})$ We have $T_x M \simeq \mathbb{R}^n$.

Heres'an example for the 2-sphere (Wikipedia):



A (differentiable) map $f: X \to Y$ of manifolds induces a linear map $T_x f: T_x X \longrightarrow T_{f(x)} Y$ $X \longmapsto f \circ Y$ (transport path to Y)

Back to the topic.

 \mathbb{R}^n , and any open subset of \mathbb{R}^n , is obviously a manifold (there is nothing to patch). Hence, $\operatorname{End}(\mathbb{R}^n) \cong \operatorname{Mot}_n(\mathbb{R}) \cong \mathbb{R}^{n^2}$

is a manifold (of dimension n²). The determinant map

$$\begin{array}{cccc}
\text{det} & Mat_n(R) \longrightarrow R \\
& A \longmapsto \text{det} A
\end{array}$$

is given by polynomials in the entries (coordinates) \Longrightarrow this is a continuous map, hence $GL_n(R) = det^{-1}(R \setminus SOS)$

$$aL_n(R) = det (R(10))$$

$$an open subset of R$$

is an open value of $Mat_n(R) \implies it$ is a (sub-) manifold and $T_x \left(GL_n(R)\right) \stackrel{naturally}{\simeq} Mat_n(R)$ $R^{n^2} \cong R^{n^2}$

Now, what about $SD(3) = \{A \in Mat_3(R) \mid A \cdot A^t = I, det(A) = 1\}$?

related to implicit function theorem.

<u>Preimage Theorem</u>: Let $f: X \rightarrow Y$ be a differentiable map of manifolds and let $y \in Y$ be a <u>regular value</u> (meaning that $T_x f: T_x X \rightarrow T_y$ is surjective $\forall x \in f'(y)$). Then $f'(y) \in X$ is a submanifold.

Apply this to the map

$$f: Mat_n(\mathcal{R}) \longrightarrow Mat_n(\mathcal{R})$$
$$A \longmapsto A \cdot A^t$$

This map is differentiable and one can show that $I \in Mat_n(R)$ is a regular value. => $O(n) = \{A \in Mat_n(R) \mid A \cdot A^t = I \} = \int^{-1} (Id)$

is a submanifold.

Note
$$A \cdot A^{\pm} - T \Rightarrow det(A)^2 = 1 \Rightarrow det A = \pm 1$$
. Hence,

$$SO(n) = (det|_{O(n)})^{-1} (R_{>O})$$

$$\xrightarrow{\text{orm-subset}} dR$$
is an open subset of $O(n) \Rightarrow$ it is a manifold. It's dimension is $n^2 - \frac{n(n+1)}{2} = \frac{n(n-1)}{2}$.
So, e.g. dim $SO(3) = 3$. Hence, dim $T_{eff}(SO(3)) = 3$, i.e. we have 3 tangent directions.
 T_{idekty}
The three tangent directions are practicly the infinitesinal generators L_x, L_y, L_z of
rotation around the $x, y, z - axis$, respectively.
Now, the tangent space is in this case not just a vector space, it has
more structure. Since $SO(n)$ is an open abramifold of $O(n)$ and $O(n)$ is
a closed submanifold of $GL_n(R)$, we naturally have an entrading
 $So(n) = T_{eff}(SO(n)) = T_{eff}(O(n)) \longrightarrow T_{eff}(GL(n)) = :g(n) \cong Mat_n(R)$
(Inder this identification, $SO(n) = 5$ skew-symmetric matrice, i.e. $A^{\pm} = -A$).
For $SO(3)$ one can identify
 $L_x = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 1 & 0 \end{bmatrix}$, $L_y = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$.
The commutator of two matrices $A_1 \otimes S(Mat_n(R)$ is
 $[A_1 \otimes B] := A \otimes B - B A$ $(L_1 \otimes B) = O(A) \otimes S(A)$:
 $[L_x, L_y] = L_z$, $[L_2, L_x] = L_y$.
These commutator velations between the infinitesival generators of rotation is of furthermatic
importance.

In general: a group which is at the same time a manifold is called a <u>Lie group</u> If G is a Lie group, one can always define a bracket [:,.] on the tangent space Lie (G) := T₁₄G satisfying similar properties as the commutator. Such a structure is called a <u>Lie algebra</u>.

The Lie algebra of a Lie group is an <u>"infinitesimal residue</u>" of the Lie group. It is so important because it is much simpler (vector space + extra structure) but still sees a (at of G. For example, if g:G→GL(V) is a representation of a lie group, then its differential gives a linear map T_Hg: Lie(G) = T_HG→T_HGL(V) = gl(V) Satisfying (T_Hg)([A,B]) = [T_H(g)(A), T_H(g)(B)], i.e. T_Hg is a representation of the Lie algebra Lie(G)! There is also a partial convose to this. Lie algebras and their representations are thus a key text in studying Lie groups and their representations. This is the motivation for this cause. A wice feature of Lie theory is that often things are constrolled by discrete combinational date ~> allows explicit calculations and computer experiments. Curious². Then follow the course! We will only be concerned with Lie algebras, no manifolds, so <u>don't worry</u> if you didn't undestand all details.